6.641 Electromagnetic Fields, Forces, and Motion Spring 2009

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6.641, Electromagnetic Fields, Forces, and Motion Prof. Markus Zahn

Lecture 4: The Scalar Electric Potential and the Coulomb Superposition Integral

I. Quasistatics

Electroquasistatics (EQS)

$$\nabla \times \overline{E} = -\frac{\partial}{\partial t} \left(\mu_0 \overline{H} \right) \approx 0$$

$$\nabla \bullet \left(\epsilon_{_0} \overline{\mathsf{E}}\right) = \rho$$

$$\nabla\times\overline{H}=\bar{J}+\frac{\partial}{\partial t}\Big(\epsilon_{_{0}}\bar{E}\Big)$$

$$\nabla \bullet \bar{J} + \frac{\partial \rho}{\partial t} = 0$$

Magnetquasistatics (MQS)

$$\nabla\times\overline{E}=-\frac{\partial}{\partial t}\Big(\mu_0\overline{H}\Big)$$

$$\nabla \times \overline{H} = \overline{J} + \frac{\partial}{\partial t} \left(z_0 \overline{E} \right)$$

$$\nabla \cdot \left(\mu_0 \overline{H}\right) = 0$$

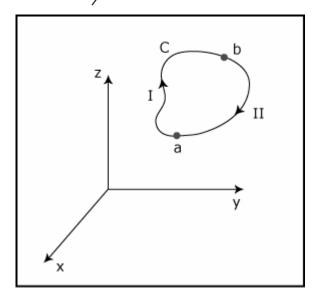
$$\nabla \cdot \bar{J} = 0$$

$$\nabla \cdot (\epsilon_0 \overline{\mathsf{E}}) = \rho$$

II. Irrotational EQS Electric Field

1. Conservative Electric Field

$$\oint_{C} \overline{E} \cdot \overline{ds} = -\frac{d}{dt} \int_{S} \mu_{0} \overline{H} \cdot \overline{da} \approx 0$$



$$\oint_{C} \overline{E} \cdot \overline{ds} = \int_{a}^{b} \overline{E} \cdot \overline{ds} + \int_{b}^{a} \overline{E} \cdot \overline{ds} = 0 \Rightarrow \int_{a}^{b} \overline{E} \cdot \overline{ds} = \int_{a}^{b} \overline{E} \cdot \overline{ds}$$

$$III$$
Electromotive Force (EMF)

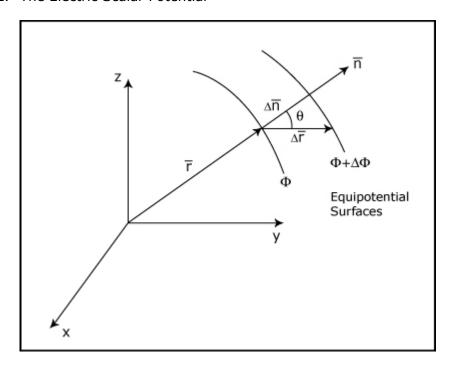
EMF between 2 points (a, b) independent of path \bar{E} field is conservative

$$\Phi\left(\bar{r}\right) \quad -\Phi\left(\bar{r}_{\mathsf{ref}}\right) = \int\limits_{\bar{r}}^{\bar{r}_{\mathsf{ref}}} \bar{E} \bullet \overline{ds}$$

electric potential

$$\int\limits_{a}^{b}\overline{E}\bullet\overline{ds}=\int\limits_{a}^{\bar{r}_{ref}}\overline{E}\bullet\overline{ds}+\int\limits_{\bar{r}_{ref}}^{b}\overline{E}\bullet\overline{ds}=\Phi\left(a\right)-\Phi\left(\bar{r}_{_{ref}}\right)+\Phi\left(\bar{r}_{_{ref}}\right)-\Phi\left(b\right)=\Phi\left(a\right)-\Phi\left(b\right)$$

2. The Electric Scalar Potential



$$\bar{r} = x \bar{i}_x + y \bar{i}_y + z \bar{i}_z$$

$$\bar{\Delta r} = \Delta x \, \bar{i}_x + \Delta y \, \bar{i}_y + \Delta z \, \bar{i}_z$$

$$\Delta n = \Delta r \cos \theta$$

$$\begin{split} \Delta \Phi &= \Phi \left(x + \Delta x, y + \Delta y, z + \Delta z \right) - \Phi \left(x, y, z \right) \\ &= \Phi \left(x, y, z \right) + \frac{\partial \Phi}{\partial x} \Delta x + \frac{\partial \Phi}{\partial y} \Delta y + \frac{\partial \Phi}{\partial z} \Delta z - \Phi \left(x, y, z \right) \\ &= \frac{\partial \Phi}{\partial x} \Delta x + \frac{\partial \Phi}{\partial y} \Delta y + \frac{\partial \Phi}{\partial z} \Delta z \\ &= \underbrace{\left[\frac{\partial \Phi}{\partial x} \bar{i}_x + \frac{\partial \Phi}{\partial y} \bar{i}_y + \frac{\partial \Phi}{\partial z} \bar{i}_z \right]}_{\text{grad } \Phi} \cdot \Delta \bar{r} \end{split}$$

$$\nabla = \bar{i}_{x} \frac{\partial}{\partial x} + \bar{i}_{y} \frac{\partial}{\partial y} + \bar{i}_{z} \frac{\partial}{\partial z}$$

$$\text{grad } \Phi = \nabla \Phi = \bar{i}_x \frac{\partial \Phi}{\partial x} + \bar{i}_y \frac{\partial \Phi}{\partial y} + \bar{i}_z \frac{\partial \Phi}{\partial z}$$

$$\int\limits_{\bar{r}}^{\bar{r}+\Delta\bar{r}} \overline{E} \bullet \overline{ds} = \Phi\left(\bar{r}\right) - \Phi\left(\bar{r}+\Delta\bar{r}\right) = -\Delta\Phi = -\nabla\Phi \bullet \Delta\bar{r} = \overline{E} \bullet \Delta\bar{r}$$

$$\overline{\mathsf{E}} = -\nabla \Phi$$

$$\Delta \Phi = \frac{\Delta \Phi}{\Delta n} \Delta r \cos \theta = \frac{\Delta \Phi}{\Delta n} - \frac{1}{\Delta r} = \nabla \Phi \cdot \Delta r$$

$$\nabla \Phi = \frac{\Delta \Phi}{\Delta n} = \frac{\partial \Phi}{\partial n} = \frac{\partial \Phi}{\partial$$

The gradient is in the direction perpendicular to the equipotential surfaces.

III. Vector Identity

$$\nabla \times \bar{\mathsf{E}} = \mathsf{0}$$

$$\bar{\mathsf{E}} = -\nabla \Phi$$

$$\nabla \times \left(\nabla \Phi \right) = 0$$

IV. Sample Problem

$$\Phi(x,y) = \frac{V_0 xy}{a^2}$$
 (Equipotential lines hyperbolas: xy=constant)

$$\begin{split} \overline{E} &= -\nabla \Phi = - \left[\frac{\partial \Phi}{\partial x} \, \overline{i}_x + \frac{\partial \Phi}{\partial y} \, \overline{i}_y \, \right] \\ &= \frac{-V_0}{a^2} \left(y \, \overline{i}_x + x \, \overline{i}_y \, \right) \end{split}$$

Electric Field Lines [lines tangent to electric field]

$$\frac{dy}{dx} = \frac{E_y}{E_x} = \frac{x}{y} \Rightarrow ydy = xdx$$

$$\frac{y^2}{2} = \frac{x^2}{2} + C$$

$$y^2 - x^2 = y_0^2 - x_0^2$$
 [lines pass through point (x_0, y_0)] (hyperbolas orthogonal to xy)

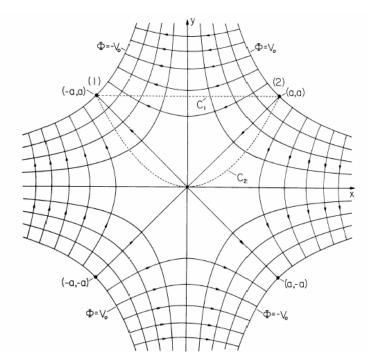


Figure 4.1.3 Cross-sectional view of surfaces of constant potential for two-dimensional potential given by (18).

Courtesy of Hermann A. Haus and James R. Melcher. Used with permission.

V. Poisson's Equation

$$\nabla \cdot \overline{\mathsf{E}} = \nabla \cdot \left(-\nabla \Phi \right) = \rho / \epsilon_0 \Rightarrow \nabla^2 \Phi = -\rho / \epsilon_0$$

$$\nabla^2 \Phi = \nabla \cdot \left(\nabla \Phi \right) = \left[\overline{\mathsf{i}}_{\mathsf{x}} \frac{\partial}{\partial \mathsf{x}} + \overline{\mathsf{i}}_{\mathsf{y}} \frac{\partial}{\partial \mathsf{y}} + \overline{\mathsf{i}}_{\mathsf{z}} \frac{\partial}{\partial \mathsf{z}} \right] \cdot \left[\frac{\partial \Phi}{\partial \mathsf{x}} \overline{\mathsf{i}}_{\mathsf{x}} + \frac{\partial \Phi}{\partial \mathsf{y}} \overline{\mathsf{i}}_{\mathsf{y}} + \frac{\partial \Phi}{\partial \mathsf{z}} \overline{\mathsf{i}}_{\mathsf{z}} \right]$$

$$= \frac{\partial^2 \Phi}{\partial \mathsf{x}^2} + \frac{\partial^2 \Phi}{\partial \mathsf{y}^2} + \frac{\partial^2 \Phi}{\partial \mathsf{z}^2}$$

VI. Coulomb Superposition Integral

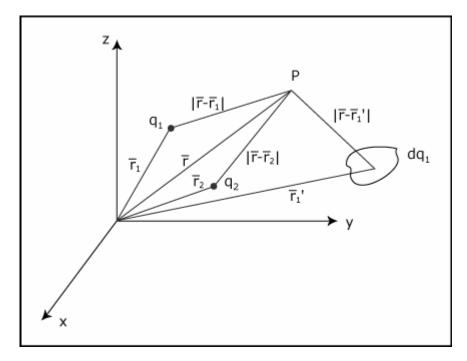
1. Point Charge

$$E_r = -\frac{\partial \Phi}{\partial r} = \frac{q}{4\pi\epsilon_0 r^2} \Rightarrow \Phi = \frac{q}{4\pi\epsilon_0 r} + C$$

Take reference $\Phi\left(r \to \infty\right) = 0 \Rightarrow C = 0$

$$\Phi = \frac{q}{4\pi\epsilon_0 r}$$

2. Superposition of Charges



$$d\Phi_{T}\left(P\right) = \frac{1}{4\pi\epsilon_{0}} \left[\frac{q_{1}}{\left|\bar{r} - \bar{r}_{1}\right|} + \frac{q_{2}}{\left|\bar{r} - \bar{r}_{2}\right|} + \dots \right. \left. \frac{dq_{1}}{\left|\bar{r} - \bar{r}_{1}\right|} + \frac{dq_{2}}{\left|\bar{r} - \bar{r}_{2}\right|} + \dots \right]$$

$$\Phi_{T}\left(P\right) = \frac{1}{4\pi\epsilon_{0}} \left[\sum_{n=1}^{N} \frac{q_{n}}{\left|\bar{r} - \bar{r}_{n}\right|} + \int\limits_{\substack{\text{all line,}\\ \text{surface, and}\\ \text{volume charges}}} \frac{dq}{\left|\bar{r} - \bar{r'}\right|} \right]$$

$$=\frac{1}{4\pi\epsilon_0}\left[\sum_{n=1}^N\frac{q_n}{\left|\bar{r}-\bar{r}_n\right|}+\int\limits_L\frac{\lambda\left(\bar{r}'\right)dl'}{\left|\bar{r}-\bar{r}'\right|}+\int\limits_S\frac{\sigma_s\left(\bar{r}'\right)da'}{\left|\bar{r}-\bar{r}'\right|}+\int\limits_V\frac{\rho\left(\bar{r}'\right)dV'}{\left|\bar{r}-\bar{r}'\right|}\right]$$

Short-hand notation

$$\Phi\left(\bar{r}\right) = \int_{V} \frac{\rho\left(\bar{r}'\right)dV'}{4\pi\epsilon_{0}\left|\bar{r}-\bar{r}'\right|}$$